

## Exam — Session 1

*Duration: 2h.*

*Documents, cell phones, computers, tablets, pocket calculators, etc., are not allowed.  
 The text contains 4 pages in total, and the 2 exercices are independent from each other.*

### 1 One-dimensional gas of hard rods (Tonks gas)

Let us consider a simple one-dimensional model for a classical fluid, composed of  $N$  hard rods of length  $\ell$  and mass  $m$  (see Fig. 1). The rods are confined along a one-dimensional space of length  $L$  and are maintained at a fixed temperature  $T$  (canonical ensemble). The rods interact through a two-body hard-wall potential  $V(x)$ . We denote by  $x_i$  the position of the middle of the  $i$ th rod ( $i = 1, \dots, N$ ).

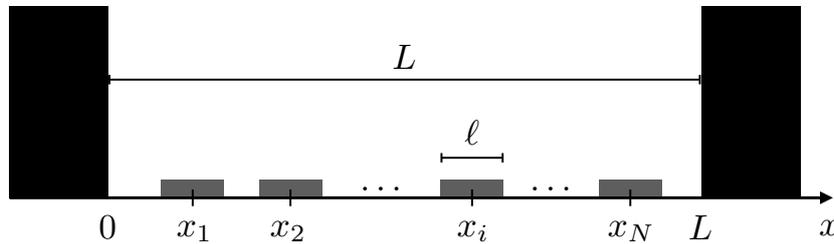


Figure 1: Sketch of a one-dimensional gas of  $N$  hard rods confined along a one-dimensional space of length  $L$  by two walls (in black).

- (a) Are the rods distinguishable or indistinguishable? Very briefly justify your answer.
- (b) Show that in the (semi-)classical, dilute limit, the canonical partition function of the system takes the form

$$Z = \left( \frac{2\pi m k_B T}{h^2} \right)^{N/2} I_N(L),$$

with  $k_B$  and  $h$  the Boltzmann and Planck constants, respectively, and where the integral  $I_N(L)$  is defined as

$$I_N(L) = \int_{\ell/2}^{L-(N-1)\ell-\ell/2} dx_1 \dots \int_{x_{i-1}+\ell}^{L-(N-i)\ell-\ell/2} dx_i \dots \int_{x_{N-1}+\ell}^{L-\ell/2} dx_N. \quad (1.1)$$

We remind the reader that  $\int_{-\infty}^{+\infty} du \exp(-u^2) = \pi^{1/2}$ .

- (c) By performing the change of variables

$$y_i = x_i + (N-i)\ell + \frac{\ell}{2} \quad (i = 1, \dots, N)$$

in Eq. (1.1), show that

$$I_N(L) = \frac{(L - N\ell)^N}{N!}.$$

- (d) Deduce from your answers to the previous questions the equation of state of the system, *i.e.*, the pressure  $P$  as a function of the parameters of the problem. Sketch the isothermal curves  $P$  as a function of the system size  $L$  for various temperatures. Does the system showcase a phase transition?

(e) The virial expansion of the equation of state is given by

$$\beta P = \sum_{n=1}^{\infty} B_n(T) \rho^n,$$

where  $\beta = 1/k_B T$ ,  $\rho = N/L$  is the density, and

$$B_2(T) = \frac{1}{2} \int dx \left[ 1 - e^{-\beta V(x)} \right].$$

Calculate in two different ways the second virial coefficient  $B_2(T)$  and check the consistency of your results.

(f) For the one-dimensional system at hand, the isothermal compressibility is defined as

$$\chi_T = -\frac{1}{L} \left( \frac{\partial L}{\partial P} \right)_T.$$

Calculate  $\chi_T$  and comment on your result.

## 2 The Blume–Capel model

The Blume–Capel model describes a magnetic material with some nonmagnetic vacancies. Let us consider a lattice [we denote by  $N (\gg 1)$  the number of lattice sites and by  $z$  the number of nearest neighbors] of spins  $S_i$  that can take the values  $-1$ ,  $0$ , and  $+1$ . A spin  $0$  corresponds to a vacancy (nonmagnetic impurity or empty site) and spins  $+1$  or  $-1$  correspond to the two different orientations of the magnetic species. We assume that the Hamiltonian of the system in presence of a homogeneous magnetic field  $h$  is given by

$$\mathcal{H} = -J \sum_{\langle i,j \rangle} S_i S_j + \Delta \sum_{i=1}^N S_i^2 - h \sum_{i=1}^N S_i, \quad (2.1)$$

where  $J > 0$  is the exchange interaction and where  $\Delta$  is a constant that can be either negative or positive. In the Hamiltonian above,  $\langle i, j \rangle$  denotes a summation over nearest neighbors.

### 2.1 General discussion

- Justify that  $-\Delta$  is the energy of creation of a vacancy. In which case ( $\Delta > 0$  or  $\Delta < 0$ ) is it favorable to create a vacancy?
- At  $T = 0$  and  $h = 0$ , calculate the energy of the system in the three different states  $S_i = +1$ ,  $S_i = -1$ , and  $S_i = 0$  ( $\forall i = 1, \dots, N$ ). Which state(s) is (are) selected at  $T = 0$ , depending on the value of  $\Delta$  with respect to  $zJ/2$ ?
- Which limit of  $\Delta$  corresponds to the usual two-state Ising model?

### 2.2 Mean-field approximation

We now aim at performing a mean-field approximation (MFA). We write  $S_i = m + \delta S_i$ , where  $m = \langle S_i \rangle$  is the average magnetization and  $\delta S_i$  the fluctuations of the spin  $S_i$  around  $m$ .

- Define the spin-spin correlation function  $C_{ij}$ . What is the value of  $C_{ij}$  in the MFA?
- Show that within the MFA, it is possible to write the Hamiltonian (2.1) as

$$\mathcal{H} \simeq \frac{1}{2} N z J m^2 - (h + z J m) \sum_{i=1}^N S_i + \Delta \sum_{i=1}^N S_i^2.$$

- (c) Calculate the free energy  $F$  within the MFA.
- (d) Demonstrate that the average value  $m = \langle S_i \rangle$  is given by the expression

$$m = -\frac{1}{N} \frac{\partial F}{\partial h}.$$

Deduce that, within the MFA, the magnetization obeys the self-consistent equation (SCE)

$$m = \frac{2 \sinh(\beta[h + zJm])}{\exp(\beta\Delta) + 2 \cosh(\beta[h + zJm])}.$$

**From now on, we consider the case of vanishing magnetic field,  $h = 0$ .**

- (e) In the case  $\Delta \rightarrow -\infty$ , discuss the solutions of the SCE.
- (f) In the general case, show that  $m = 0$  is a solution of the SCE.
- (g) We now aim at discussing graphically the solutions of the SCE. We define  $t = k_B T / zJ$  and  $\delta = \Delta / zJ$ .

- (i) Express the SCE in term of the function

$$f(m) = \frac{2 \sinh(m/t)}{\exp(\delta/t) + 2 \cosh(m/t)}.$$

- (ii) What is the value of  $f(0)$ ?
- (iii) What are the limits of  $f(m)$  when  $m \rightarrow \pm\infty$ ?
- (iv) Calculate

$$\left. \frac{df}{dm} \right|_{m=0}$$

and discuss graphically the number of solutions to the SCE. Show that there is a critical reduced temperature  $t_c$  defined by the equation

$$t_c = \frac{2}{2 + \exp(\delta/t_c)}.$$

- (v) On next page in Fig. 2 (colored lines) is plotted the function

$$g(t, \delta) = \frac{2}{2 + \exp(\delta/t)} \quad (2.2)$$

as a function of  $t$  for different values  $\delta_i$  of  $\delta$ . Which  $\delta_i$ 's are positive and which of them are negative? Sort by ascending order the  $\delta_i$ 's.

- (vi) Plot the curve  $g(t, \delta)$  for the value of  $\delta$  corresponding to the Ising model and give the corresponding  $t_c$ .
- (vii) Using your previous discussion and question 2.1(b), sketch the general behavior of  $t_c$  as a function of  $\delta$ .

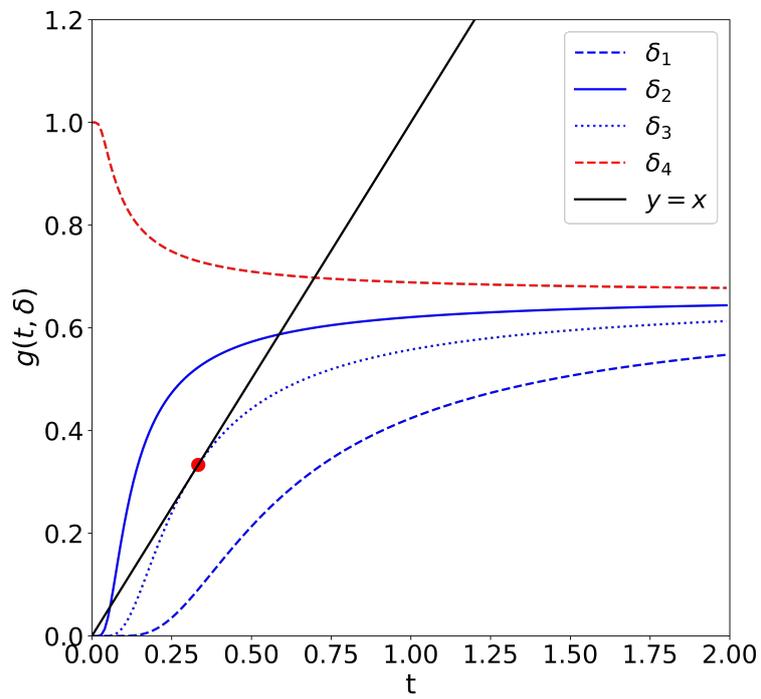


Figure 2: Colored lines: Plot of  $g(t, \delta)$  as defined in Eq. (2.2) as a function of  $t$  for different values  $\delta_i$  of  $\delta$ . Black solid line:  $t$ .