Exam - Session 1

Duration: 2h.

Documents, cell phones, computers, tablets, pocket calculators, etc., are not allowed. The text contains 3 pages in total, and the 2 exercises are independent from one another.

1 Charge carrier transport in a doped semiconductor

We consider a semiconductor with equilibrium electron and hole concentrations n_0 and p_0 , respectively. In the following, we shall address carrier transport phenomena and the *non-equilibrium* electron and hole concentrations are denoted $n(\mathbf{r},t)$ and $p(\mathbf{r},t)$, respectively. The elementary charge is denoted e(>0). We introduce the electron and hole mobilities, μ_n and μ_p , respectively. The diffusion coefficient of electrons and holes are denoted D_n and D_p , respectively.

1.1 Drift and diffusion currents

- (a) In the presence of an external electric field **E**, express the electron and hole *drift* current densities, denoted \mathbf{J}_{n}^{drift} and \mathbf{J}_{p}^{drift} , respectively. For a given orientation of **E**, draw the electron and hole fluxes as well as the electron and hole current densities. Connect the electron (hole) mobilities to the electron (hole) conductivities $\sigma_{n(p)}$.
- (b) In the presence of charge carrier concentration gradients, we define the electron and hole *diffusion* current densities $\mathbf{J}_{n}^{\text{diff}}$ and $\mathbf{J}_{p}^{\text{diff}}$, respectively. Give their expressions as a function of $n(\mathbf{r},t)$ $[p(\mathbf{r},t)]$ and D_{n} $[D_{p}]$ and, as in the previous question, qualitatively draw, both for electrons and holes, the charge carrier concentration gradient, the associated particle flux, and the associated diffusion current density.
- (c) Write the *total* current density \mathbf{J}^{tot} .

1.2 Doped semiconductor

We now consider an homogeneously doped semiconductor with an electronic gap \mathcal{E}_{G} and we focus on the case of an electron-doped (*i.e.*, *n*-doped) system, where doping arises from an impurity band, with donor concentration N_{d} , situated at an energy $\delta \mathcal{E} \ll \mathcal{E}_{G}$ below the conduction band minimum.

- (a) Give an example of an *n*-doped semiconductor. Explain qualitatively (no calculations are requested) how n_0 can be determined in the following three temperature regimes: $k_{\rm B}T \ll \delta \mathcal{E}$ (frozen regime); $k_{\rm B}T \gtrsim \delta \mathcal{E}$ (saturation regime); $k_{\rm B}T \gg \delta \mathcal{E}$ (intrinsic regime). The temperature is denoted T and $k_{\rm B}$ is the Boltzmann constant.
- (b) Draw qualitatively $\ln(n_0)$ as a function of 1/T.
- (c) In the following, we shall consider the saturation regime. Express the charge neutrality condition. Knowing that the *intrinsic* charge concentration at room temperature is typically $n_{\rm i} \sim 10^{16} \text{ m}^{-3}$ and that $N_{\rm d} \sim 10^{20} \text{ m}^{-3}$, give a numerical estimate of p_0 . Comment on this result and justify why holes are called "minority carriers" in an *n*-doped semiconductor.

1.3 Continuity equation for minority carriers in an *n*-doped semiconductor

In the following we consider, for simplicity, a one-dimensional (1d) problem along the x axis.

- (a) Write the continuity equation that connects the hole concentration and $J_{\rm p}^{\rm tot}$ in the absence of generation and recombination of charge carriers.
- (b) We now consider generation and recombination of electron-hole pairs. How can electronhole pairs be generated and then recombine in a semiconductor? Justify precisely that the difference between the electron-hole pair generation and recombination rates can be written as $\beta (n_0 p_0 - np)$, where the (x, t) dependence of n and p is implicitly considered. What is the physical meaning and the dimension of β ?
- (c) Still in the case of an *n*-doped semiconductor, demonstrate that one can approximate the difference between the generation and recombination rates as $(p_0 p)/\tau_p$. What is the physical meaning and the dimension of τ_p ?
- (d) Demonstrate that the continuity equation for holes writes

$$\frac{\partial p}{\partial t} = D_{\rm p} \frac{\partial^2 p}{\partial x^2} - \mu_{\rm p} \frac{\partial (pE)}{\partial x} + \frac{p_0 - p}{\tau_{\rm p}}.$$
(1.1)

1.4 Diffusion of minority carriers in a 1d *n*-doped semiconductor

We consider a 1d *n*-doped semiconductor in which a constant excess of carriers (electrons and holes) $\Delta p_0 = \Delta n_0$ is maintained in the region $x \leq 0$, whereas no excess carriers are injected in the region x > 0, such that $\lim_{x\to\infty} p(x) = p_0$. For now, we neglect the internal electric field E(x) that may arise within the semiconductor. This approximation will be relaxed in question 1.4(d).

- (a) Using Eq. (1.1) and the boundary conditions, express $\Delta p(x) = p(x) p_0$ for any x and draw $\Delta p(x)$. You may introduce $L_p = \sqrt{D_p \tau_p}$. What is the physical meaning of L_p ?
- (b) Assuming local quasi-neutrality, one can write $\Delta n(x) = n(x) n_0 \approx \Delta p(x)$. Express $J_{\rm p}^{\rm diff}(x)$ and $J_{\rm n}^{\rm diff}(x)$.
- (c) What must be the value of the total current density $J^{\text{tot}}(x)$?
- (d) Using this condition, express the electric field E(x) as a function of $L_{\rm p}$, $D_{\rm n}$, $D_{\rm p}$, $\sigma_{\rm n}$ and $\sigma_{\rm p}$. Comment on this result.

We now consider that a constant excess of holes $\Delta p_{\rm l}$ is maintained for $x \leq 0$, whereas a similar deficit of holes $\Delta p_{\rm r}$ is maintained for $x \geq w$ (w > 0). In the following, we neglect the internal electric field.

- (e) Solve again the continuity equation (1.1) and, using the boundary conditions, express $\Delta p(x)$ in the region $0 \leq x \leq w$. Give simplified expressions for $\Delta p(x)$ in the two limiting cases $w \gg L_p$ (long device configuration) and $L_p \gg w$ (short device configuration).
- (f) Plot $\Delta p(x)$ in the two configurations introduced above.
- (g) Determine the simplified expression of $J_{\rm p}^{\rm diff}(x)$ in the short and long device configuration, respectively. Comment on the differences between these two expressions.

2 Paramagnetism of degenerate and nondegenerate free electrons

We consider a three-dimensional (3d) gas of $N \gg 1$ noninteracting free electrons of mass m and spin s = 1/2 placed in a magnetic field B directed along the z axis. Each electron has a magnetic moment of modulus equal to the Bohr magneton $\mu_{\rm B}$ with two possible orientations with respect to the applied magnetic field. We note T the temperature, $\beta = 1/k_{\rm B}T$, and μ the electron chemical potential.

- (a) Show that the density of states g(E) per unit volume for free electrons in a 3d system in the *absence* of a magnetic field as a function of the energy E is given by $g(E) = K\sqrt{E}$, where K is a constant to be determined.
- (b) We apply a weak magnetic field such that $\mu_{\rm B}B \ll \mu$. Considering only the effect of the additional Zeeman term on the spin degree of freedom, give the new dispersion relation (with respect to the B = 0 case) and the densities of states $g_{\pm}(E)$ for spins parallel (+) and antiparallel (-) to the magnetic field B, respectively.
- (c) In the degenerate case of $\beta \mu \gg 1$, indicate graphically the occupation as a function of energy f(E) of electronic states at T = 0 K and at T > 0 K.
- (d) Calculate the number of electron spins per unit volume parallel (n_+) and antiparallel (n_-) to the magnetic field.
- (e) Determine the magnetization M and the paramagnetic susceptibility χ of electrons at T = 0 K.
- (f) We now consider the nondegenerate case of low electron density, where we can approximate the Fermi–Dirac distribution by $f(E) \approx e^{-\beta(E-\mu)}$. Calculate M and χ as a function of the electron density and the temperature. Explain why the Curie law is obtained for the susceptibility of nondegenerate electrons.